

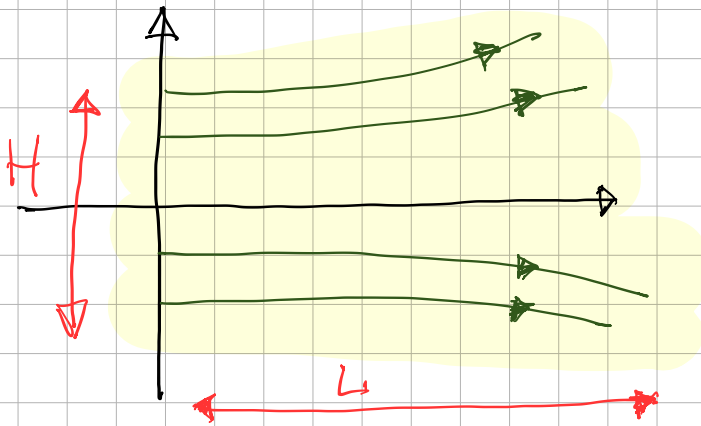
Lecture 6.

I. Parabolic equation of diffraction theory

Consider Helmholtz equation

$$\Delta \tilde{u} + k^2 \tilde{u} = 0$$

Let u be a paraxial wave beam moving along x -axis



How to describe such a beam?

$$k^{-2} \ll H \ll L$$

High frequency case

Let $\tilde{u}(x, y) = e^{ikx} U(x, y)$

amplitude function

$$\frac{\partial U}{\partial x} \sim \frac{U}{L}, \quad \frac{\partial U}{\partial y} \sim \frac{U}{H}$$

$$\left(\underbrace{\frac{\partial^2}{\partial y^2}}_{H^{-2}} + 2ik \underbrace{\frac{\partial}{\partial x}}_{KL^{-1}} + \underbrace{\frac{\partial^2}{\partial x^2}}_{L^{-2}} \right) U(x, y) = 0$$

$$\left(\frac{\partial^2}{\partial y^2} + 2ik \frac{\partial}{\partial x} \right) U(x, y) = 0$$

$$H \sim \sqrt{\frac{E}{k}}$$

a) due to $\frac{\partial}{\partial x}$ waves propagate from left to right

Consider:

$$u^{in} = e^{ikx \cos \theta} + ik y \sin \theta$$

$$\approx |\theta \ll 1| \approx e^{ikx} e^{-ikx \frac{\theta^2}{2}} + ik y \theta \approx$$

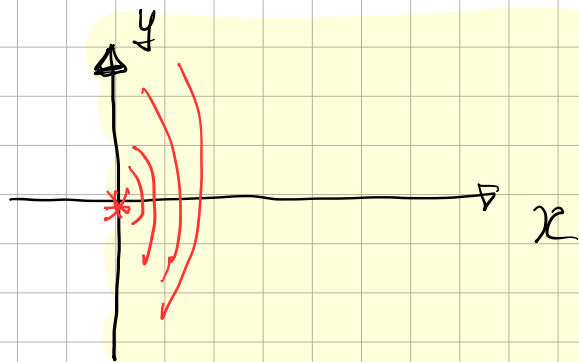
$$u^{in} = e^{-ikx \frac{\theta^2}{2}} + ik y \theta \rightarrow \text{parabolic plane wave}$$

2. Green's function of the parabolic equation

$$\left(\frac{\partial^2}{\partial y^2} + 2ik \frac{\partial}{\partial x} \right) U(x, y) = 0$$

Let's solve it for $x > 0$ with the boundary (initial) condition

$$U(0, y) = \delta(y)$$



Denote the solution

$$U(x, y) = G(x, y)$$

Let's look for the solution as plane wave decomposition

$$G(x, y) = \int_{-\infty}^{\infty} A(\xi) e^{i\xi y - i\xi^2 x / (2k)} d\xi$$

From boundary condition

$$\int_{-\infty}^{\infty} A(\xi) e^{i\xi y} d\xi = \delta(y),$$

i.e. $A(\xi) = \frac{1}{2\pi}$

$$G(x, y) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \exp \left\{ i\xi y - \frac{i\xi^2 x}{2k} \right\} d\xi =$$

$$= \int \xi y - \frac{\xi^2 x}{2k} = -\frac{x}{2k} \left(\xi - \frac{ky}{x} \right)^2 + \frac{ky^2}{2x} \quad \left(\xi = \right)$$

$$= \frac{1}{2\pi} e^{\frac{iky^2}{2x}} \int_{-\infty}^{\infty} e^{-\frac{ix}{2k} z^2} dz \quad (\text{z})$$

$$\int_{-\infty}^{\infty} e^{-ia z^2} = \sqrt{\pi/a} e^{-\frac{i\pi}{4}} a^{-\frac{1}{2}}$$

$$\Rightarrow \sqrt{\frac{k}{2\pi x}} e^{\frac{iky^2}{2x} - \frac{i\pi}{4}}$$

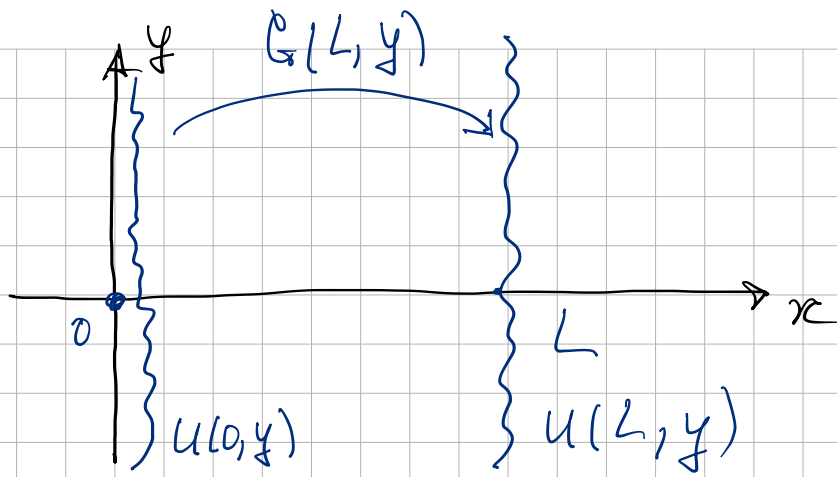
$G(x, y)$ acts as a propagator:

Let $U(0, y) = f(y)$, Then

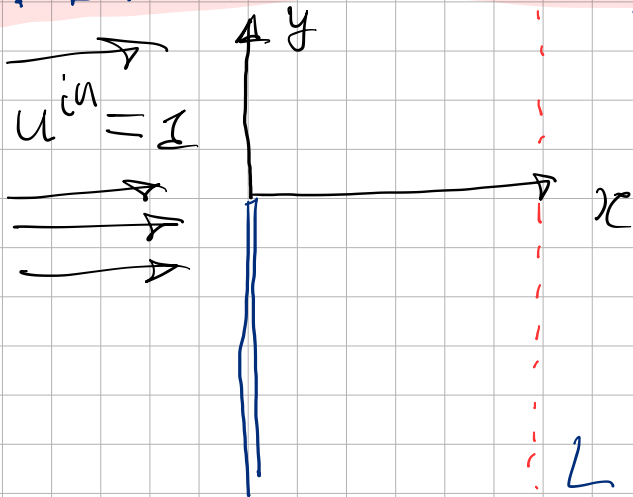
$$f(y) = \int_{-\infty}^{\infty} \delta(y - y') f(y') dy'$$

Thus

$$U(x, y) = \int_{-\infty}^{\infty} G(x, y - y') f(y') dy'$$



3. Diffraction by a half-plane (perpendicular)



$$u(L, y) = \sqrt{\frac{k}{2\pi L}} \int_0^{\infty} e^{\frac{ik(y-y')^2}{2L} - \frac{i\pi}{4}} dy'$$

can be expressed in terms of Erf

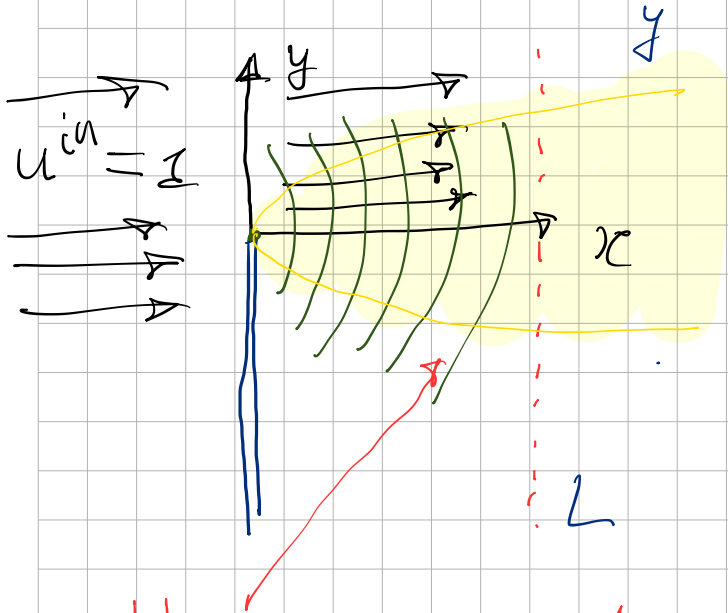
Let's estimate asymptotically

$$z = \sqrt{\frac{k}{2L}} (y' - y)$$

$$u(L, y) = \frac{e^{-\frac{i\pi}{4}}}{\sqrt{\pi}} \int_{-\sqrt{\frac{k}{2L}} y}^{\infty} e^{iz^2} dz$$

For $y > 0$ we have a stationary point $z=0$, thus

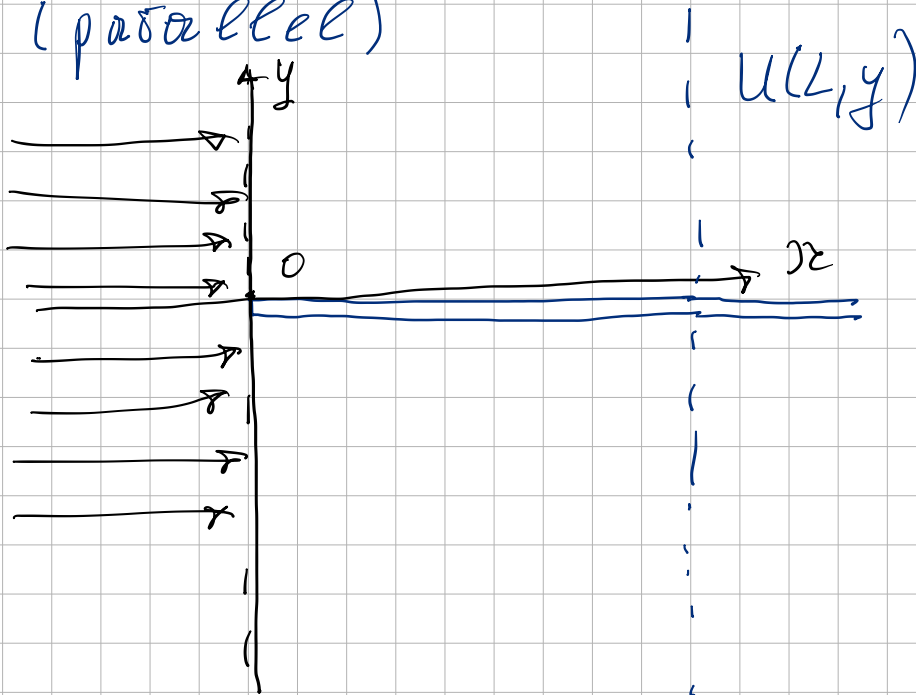
$$U(L, y) \approx \begin{cases} 1 + \frac{e^{\frac{+i\pi}{4}}}{\sqrt{8}} \frac{e^{\frac{iky^2}{2L}}}{2y\sqrt{\frac{k}{2L}}} & ; y > 0 \\ + \frac{e^{\frac{+i\pi}{4}}}{\sqrt{8}} \frac{e^{\frac{iky^2}{2L}}}{2y\sqrt{\frac{k}{2L}}} & ; y < 0 \end{cases}$$



This term can
 be obtained
 by integration
 by parts

Here the result will be close to the exact solution

4. Diffraction by a half-plane (parallel)



Similarly

$$U(L, y) \approx \frac{e^{-\frac{i\pi}{4}}}{\sqrt{\pi}} \int_{-\sqrt{\frac{k}{2L}}|y|}^{\infty} e^{iz^2} dz$$

WSON of sec comment in Lecture 7!

Remark

$$u^{sc} \approx G(x, y) S(\theta);$$

$$S(\theta) = \frac{i}{k\theta}, \text{ where } \theta = \frac{y}{x}$$

5. Directivity in the parabolic approximation

$$u^{sc}(x, y) = S(\theta) \cdot G(x, y) + O((kx)^{-3/2})$$

$\theta = \frac{y}{x}$, i.e. θ is \tan of the angle of observation

We use PEDT when $\tan \theta \sim \theta$

Let u^{sc} be known on a line



Using propagator

$$u^{sc}(x, y) \approx \int_{-\infty}^{\infty} u^{sc}(0, y') G(x, y - y') dy' = \sqrt{\frac{k}{2\pi x}} e^{-\frac{i\pi}{4}} \int_{-\infty}^{\infty} u^{sc}(0, y') e^{\frac{ik(y-y')^2}{2x}} dy' =$$

$$= \sqrt{\frac{k}{2\sqrt{8}\pi}} e^{-\frac{i\sqrt{8}}{4}} e^{\frac{ik y^2}{2\pi}} \int_{-\infty}^{\infty} u^{sc}(0, y') e^{-\frac{iky y'}{\pi} + \frac{ik(y')^2}{2\pi}} dy'$$

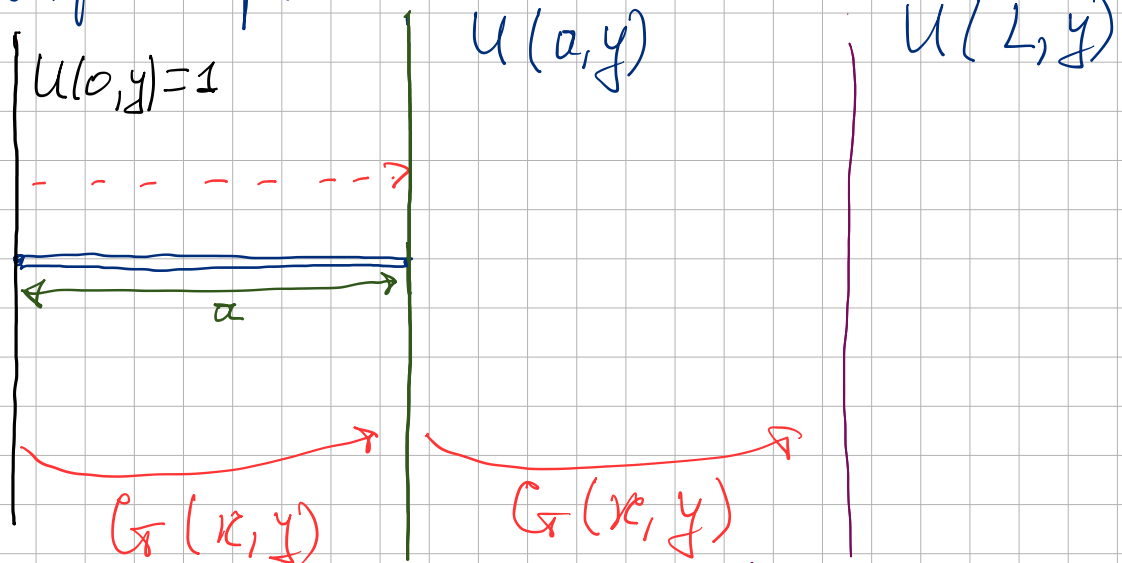
Let $x, y \rightarrow \infty$ with $\theta = \frac{y}{x} = \text{const}$

Then:

$$S(\theta) = \int_{-\infty}^{\infty} u^{sc}(0, y) e^{-iky\theta} dy$$

Homework.

Strip problem



Korolev A.I., Shanin A.V. Acoust. Journal V. 62 p. 399 (2016)

Formally $u(L, y) = \Pi_1 \Pi_2 u(0, y)$

Bonus: Asymptotic of Fresnel integral

$$F(z) = \frac{e^{-\frac{i\sqrt{8}}{4}}}{\sqrt{8}} \int_z^{\infty} 2i\tau e^{i\tau^2} \frac{d\tau}{2i\tau} ; z > 0$$

$$F(z) = \frac{e^{-\frac{i\sqrt{8}}{4}}}{\sqrt{8}} \frac{e^{i\tau^2}}{2i\tau} \Big|_z^{\infty} + \int_z^{\infty} \frac{e^{i\tau^2}}{2i\tau^2} d\tau \approx$$

$$\approx + \frac{e^{+i\frac{\sqrt{8}}{4}}}{2\sqrt{8}} e^{iz^2} \left(\frac{1}{z} + O\left(\frac{1}{z^2}\right) \right)$$