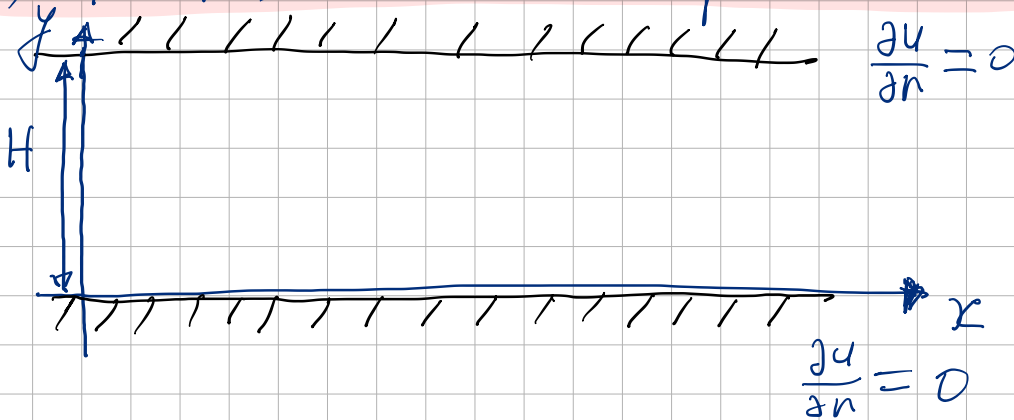


# Lecture 9. Waveguides

## 1) Modes of a plane waveguide



$$\Delta u + k^2 u = 0$$

$$u = \ell(x) p(y)$$

$$\frac{d^2 \ell}{dx^2} + k_x^2 \ell = 0$$

$$\frac{d^2 p}{dy^2} + k_y^2 p = 0$$

$$k_x^2 + k_y^2 = k^2 = \frac{\omega^2}{c^2}$$

$$a) \ell = e^{i k_x x} + e^{-i k_x x}$$

b) Sturm - Liouville (SL) problem

$$\frac{d^2 p}{dy^2} + k_y^2 p = 0 \quad \frac{dp}{dy}(0) = \frac{dp}{dy}(H) = 0$$

Find  $k_y$  for which solutions exist

$$p(y) = A \cos(k_y y) \quad k_y = \frac{\pi n}{H}$$

$$k_x^n = \sqrt{\left(\frac{\omega}{c}\right)^2 - \left(\frac{\pi n}{H}\right)^2}$$

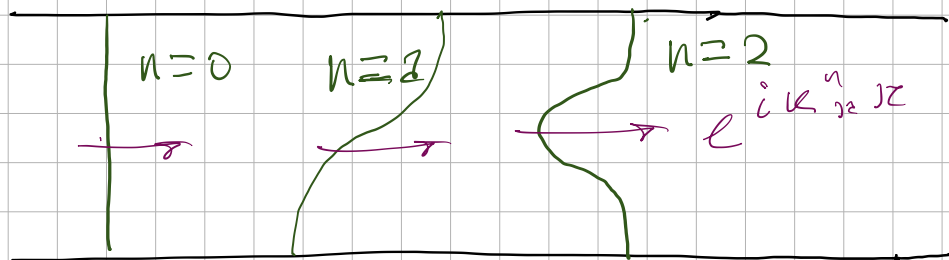
Thus

$$U_f = \sum_{n=0}^{\infty} A_n \cos\left(\frac{\pi n}{H} y\right) e^{i k_{x,n} x} =$$

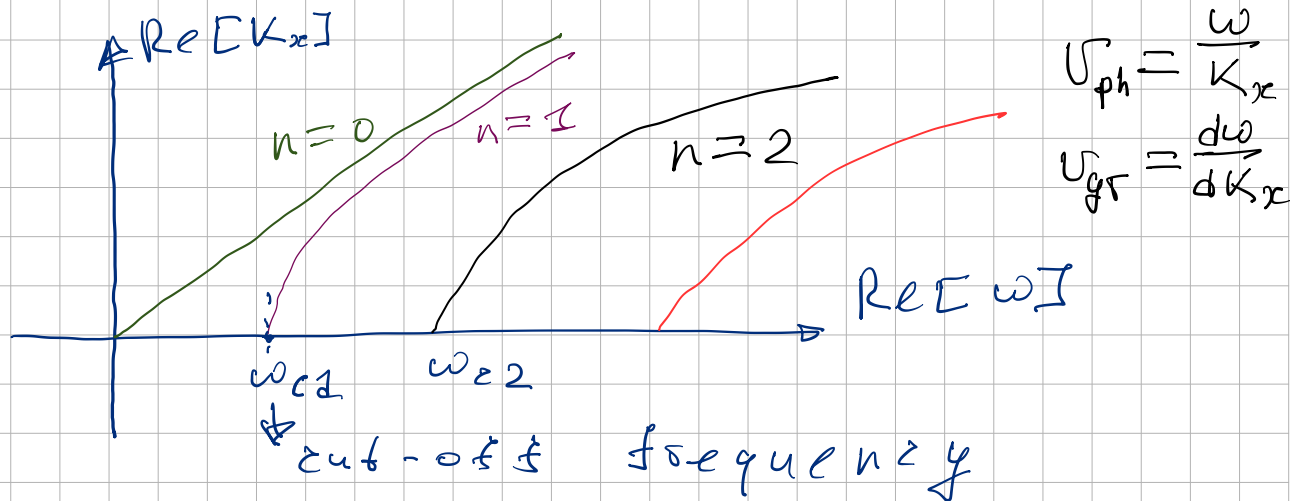
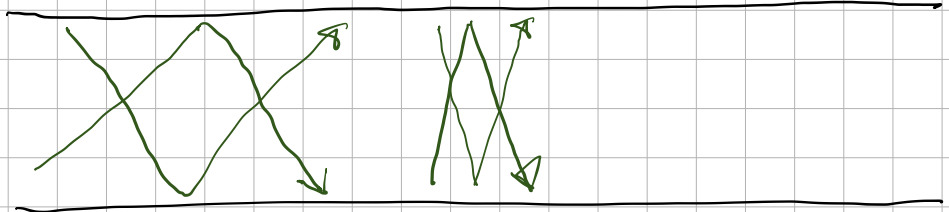
$$= \sum_{n=0}^{\infty} \tilde{A}_n \left[ e^{i k_{x,n} x} e^{-i \left(\frac{\pi n}{H}\right) y} + e^{i k_{x,n} x} e^{+i \left(\frac{\pi n}{H}\right) y} \right]$$

↓ Brillouin (partial) wave

Modes:



Partial waves



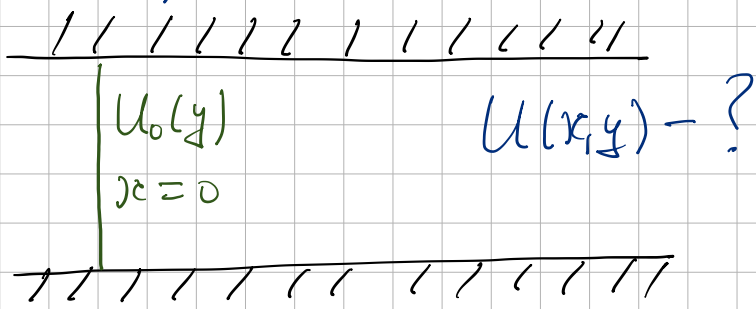
Note that modes are orthogonal:

$$\int_0^H \cos\left(\frac{\pi m}{H} y\right) \cos\left(\frac{\pi n}{H} y\right) dy = 0$$

if  $m \neq n$

Any wave field can be presented in terms of mode (Fourier) series

In deed, consider an excitation problem



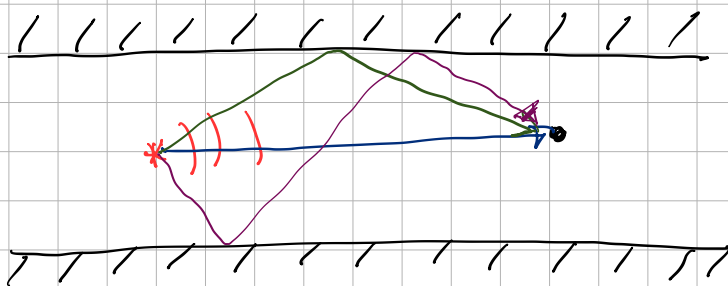
$$U_0(y) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{\pi n}{H} y\right)$$

$$A_0 = \frac{1}{H} \int_0^H U_0(y) dy$$

$$A_n = \frac{2}{H} \int_0^H U_0(y) \cos\left(\frac{\pi n}{H} y\right) dy; n > 0$$

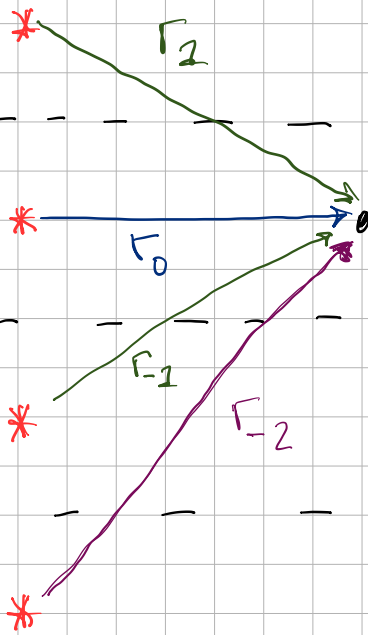
$$U(x, y) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{\pi n}{H} y\right) e^{i k_x^n x}$$

c) Green's function of a waveguide



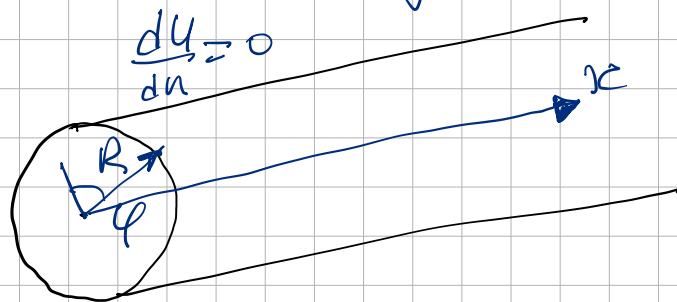
One approach is to use mode decomposition as above.

Another one is to apply method of reflections.



$$G(x, y) = \frac{-i}{4} \sum_{n=-\infty}^{\infty} H_0^2(kr_n)$$

## 2) Circular waveguide



$$\frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} + \frac{\partial^2 U}{\partial \varphi^2} + \frac{1}{r^2} \frac{\partial^2 U}{\partial \varphi^2} + k^2 U = 0$$

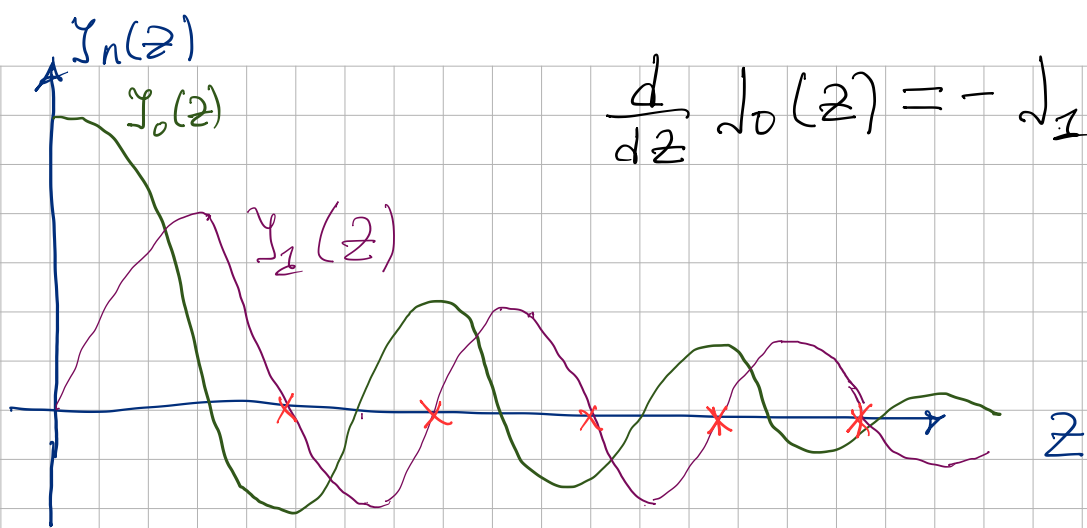
$$U = \sum_{m, n=0}^{\infty} A_{m, n} e^{i\varphi m + k_{m, n} x} F_{m, n}(r)$$

$$\frac{d^2 F}{dr^2} + \frac{1}{r} \frac{dF}{dr} + (k^2 - k_{m, n}^2 - \frac{m^2}{r^2}) F = 0$$

$$F_{m, n} = J_m(\sqrt{k^2 - k_{m, n}^2} r)$$

$$\frac{dF_{m, n}}{dr} = 0 \Rightarrow \frac{dJ_m(\sqrt{k^2 - k_{m, n}^2} r)}{dr} = 0$$

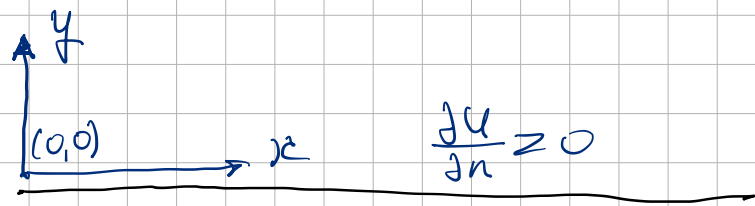
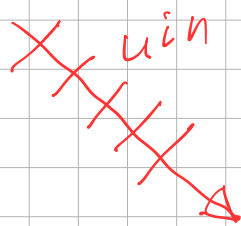
# modes are zeros of derivative



$$\frac{d}{dz} J_0(z) = -J_1(z)$$

3) Diffraction by an open-ended plane waveguide. Weinstein's problem

a) Statement



$$u = u^{sc} + u^{in}$$

$$\Delta u + k^2 u = 0$$

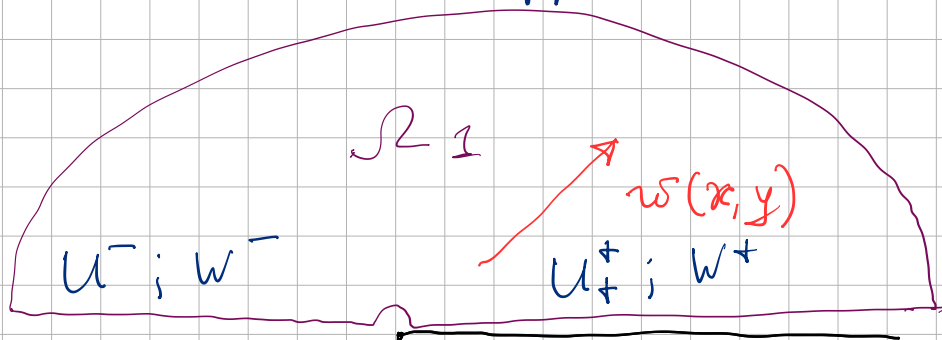
Meixner and radiation conditions

b) Derivation of Wiener-Hopf equation using Green's identity

Auxiliary function:

$$w(x, y) = e^{i\beta x + i\sqrt{k^2 - \beta^2} y}$$

Domain 1: Upper Half-plane



$$\int_{\partial \Omega_1} \left[ \frac{\partial w}{\partial n} u^{sc} - \frac{\partial u^{sc}}{\partial n} w \right] d\ell = 0$$

$$\rightarrow i\sqrt{k^2 - \zeta^2} U^- + W^- - i\sqrt{k^2 - \zeta^2} U^+ + W^+ = 0$$

↓  
Forcing

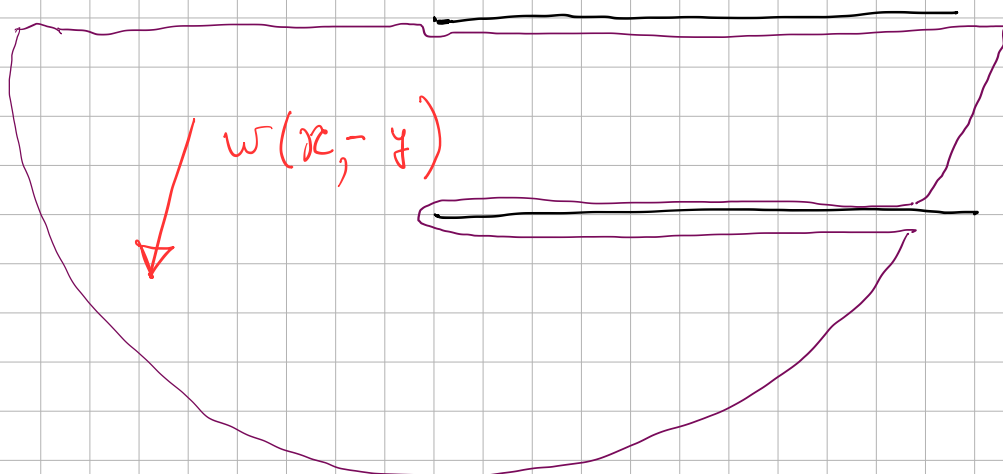
$$U^- = \int_0^{\infty} e^{i\zeta x} u^{sc}(x, 0) dx$$

$$U^+ = \int_{-\infty}^{\infty} e^{i\zeta x} u^{sc}(x, 0^+) dx$$

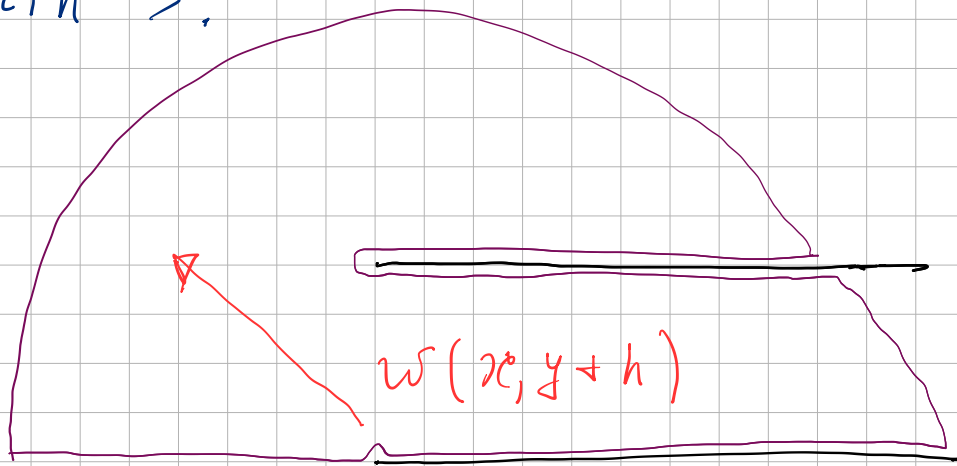
$$W^- = \int_0^{\infty} e^{i\zeta x} \frac{\partial u^{sc}}{\partial y}(x, 0) dx$$

$$W^+ = - \int_{-\infty}^{\infty} e^{i\zeta x} \frac{\partial u^{in}}{\partial y}(x, 0) dx$$

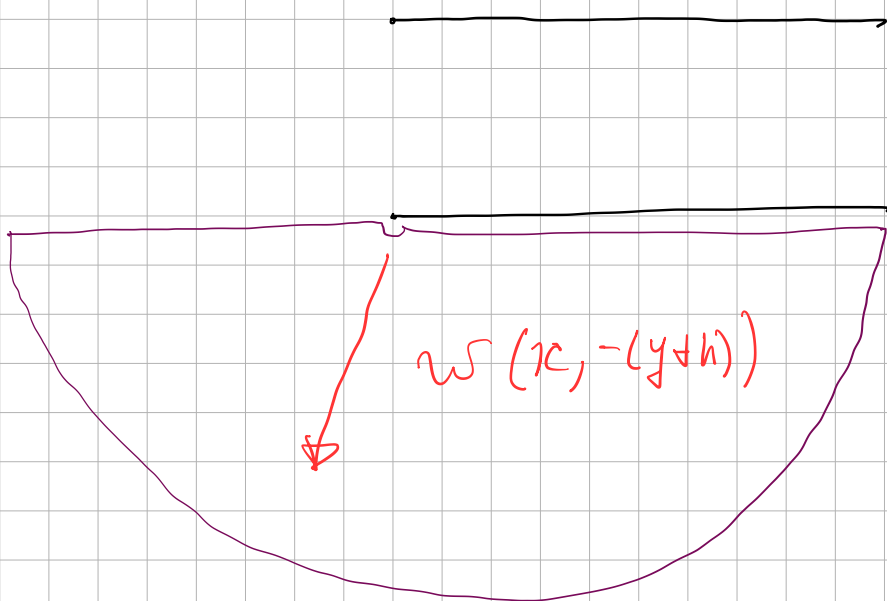
Domain 2:



Domain 3:



Domain 4:



We got 3 more functional equations, which are reduced to 2 by 2 matrix  $W-H$ :

$$\Psi^- = K \Psi^+ + F$$

$$\Psi^- = \begin{pmatrix} w^- \\ \vec{\Gamma}^- \\ \vec{\Delta}^- \end{pmatrix}; \quad \Psi^+ = \begin{pmatrix} u^+ \\ \varphi^+ \end{pmatrix}$$

$$\varphi^+ = \int_0^\infty e^{i\beta x} \{u^{sc}(x, -h^+) - u^{sc}(x, -h^-)\} dx$$

$$\vec{\Psi}^- = \int_{-\infty}^0 e^{i\zeta x} \frac{\partial u^{sc}}{\partial y}(x, -h) dx$$

$$U^+ = U_+^+ - U_-^+$$

$$K = \frac{i\zeta}{2} \begin{pmatrix} 1 & e^{i\zeta h} \\ e^{i\zeta h} & 1 \end{pmatrix}; \zeta = \sqrt{k^2 - \zeta^2}$$

$$F = \frac{k \sin \theta^{in}}{\zeta - k \cos \theta^{in}} \begin{pmatrix} 1 \\ e^{i k h \sin \theta^{in}} \end{pmatrix}$$

System can be decoupled. Indeed

$$K = -\frac{i\zeta}{4} \begin{pmatrix} -1 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1 - e^{i\zeta h} & 0 \\ 0 & 1 + e^{i\zeta h} \end{pmatrix} \begin{pmatrix} 1 & -1 \\ -1 & -1 \end{pmatrix}$$

Introduce

$$\tilde{\Psi}^- = \begin{pmatrix} 1 & -1 \\ -1 & -1 \end{pmatrix} \Psi^- \quad \tilde{\Psi}^+ = \begin{pmatrix} 1 & -1 \\ -1 & -1 \end{pmatrix} \Psi^+$$

$$\tilde{F} = \begin{pmatrix} 1 & -1 \\ -1 & -1 \end{pmatrix} F$$

$$\tilde{\Psi}^- = \frac{i\zeta}{2} \begin{pmatrix} 1 - e^{i\zeta h} & 0 \\ 0 & 1 + e^{i\zeta h} \end{pmatrix} \tilde{\Psi}^+ + \tilde{F}$$

2 independent scalar equations that are solved using scalar W-H method

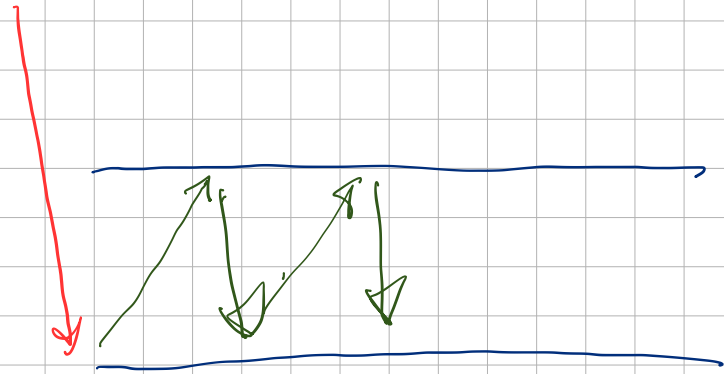
How to work: factorize  $e$

$$f(\zeta) = f_+(\zeta) f_-(\zeta)$$

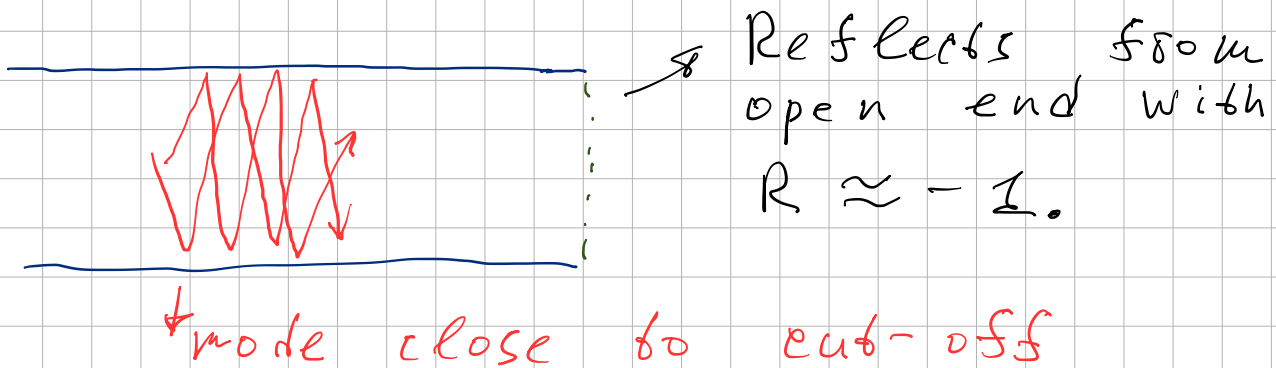
$$f(\zeta) = \sqrt{k^2 - \zeta^2} \left( 1 \pm e^{i\sqrt{k^2 - \zeta^2} h} \right)$$

# Speculations

a) Grazing wave excites waveguide



b) Vice-versa Mode close to cut-off reflects almost completely from the open end:



Home work: <sup>\*\*\*</sup> Formulate W-H problem for waveguide mode, show that almost complete reflection happens close to cut-off

Such modes form standing waves in Fabry-Perot resonators:

